

# Comments on “Observation of Strong Quantum Depletion in a Gaseous Bose-Einstein Condensate” [cond-mat/0601184]

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A recent paper by K. Xu *et al.* [1] discussed a quantum depletion in gaseous Bose-Einstein condensation. In the ground state of an interacting Bose gas, not all the particles are in the lowest energy state, because the two-body interaction mixes the ground state components with atoms in the other states. For a homogeneous system the particle depletion is known well from Bogoliubov theory as [2]  $\eta_0 = \frac{8}{3\sqrt{\pi}}\sqrt{\rho a^3}$ , where  $\rho$  is the atomic density and  $a$  is the s-wave scattering length. The authors put the data of  $^{23}\text{Na}$  into above relation and say that “For  $^{23}\text{Na}$  at a typical density of  $10^{14}\text{cm}^{-3}$ , the quantum depletion is 0.2%”. Apparently, this is not a reliable value for the inhomogeneous system which a magneto-optical trap is applied. For an interacting and inhomogeneous system the difficulty comes from the fact that we do not know the every energy eigenvalue of the system except  $\varepsilon_0$  which is obtained by a macroscopic mean field wave function so called Gross-Pitaevskii equation(GPE). However, one of the authors, W. Ketterle, already discussed the problem in an inhomogeneous system seriously and obtained an effective method for that [3]. We'll apply this method to obtain a better prediction.

We first consider a 3D system of  $N$  non-interacting Bosons confined by a harmonic potential of angular frequency  $\omega$ . When the Bose distribution function in the grand canonical ensemble at temperature  $T$  is  $f_B(\varepsilon)$  and the density of states of the isotropic 3D harmonic oscillator is  $\rho(\varepsilon)$ , the number of particles at temperature  $T$  is obtained by  $N = \int f_B(\varepsilon)\rho(\varepsilon)d\varepsilon$  [3, 4, 5, 6]. It is

$$N = N_0 + \left(\frac{k_B T}{\hbar\omega}\right)^3 g_3(z) + \frac{2}{3} \left(\frac{k_B T}{\hbar\omega}\right)^2 g_2(z) + \dots, \quad (1)$$

where  $g_n(z) = \sum_{l=1}^{\infty} z^l/l^n$  is the Bose functions.  $z = e^{\beta(\mu-\varepsilon_0)}$  is the fugacity between 0 and 1,  $\varepsilon_0 = (3/2)\hbar\omega$  is the zero-point energy of the harmonic oscillator, and  $\mu(T)$  is the chemical potential.  $N_0$  is the ground state occupation number given by  $N_0 = 1/(z^{-1} - 1)$ . The others are common notations.

For interacting system there is a shift of the zero point energy by the interaction. Then, the new ground state occupation number is written as

$$N_0^{int} = \frac{1}{e^{\beta(\varepsilon_0+U-\mu)} - 1}, \quad (2)$$

where  $U = U(a, N_0)$  is the interaction energy between atoms and may obtained from the solution of the GPE. We consider the positive scattering length only, so does

$U > 0$ . Then, the condensate fraction of the interacting system is definitely lower than that of the ideal system.

The chemical potential  $\mu(T)$  is obtained numerically by the following way. The transition temperature  $T_c$  is determined that at the onset of condensation  $N_0 \rightarrow 0$  and  $z \rightarrow 1$  from Eq. (1). It is

$$T_c = T_0 \left\{ 1 - \frac{\zeta(2)}{2\zeta(3)^{2/3}} \frac{1}{N^{1/3}} + \dots \right\}, \quad (3)$$

where  $T_0 = (\hbar\omega/k_B) \{N/\zeta(3)\}^{1/3}$  and  $\zeta$  is the Riemann-zeta function.  $T_c = 328 \sim 713\text{nK}$  for  $N = 10^5 \sim 10^6$  and  $\omega = 10^3\text{sec}^{-1}$ . The condensate fraction is obtained from Eq. (1) as a function of temperature [3, 4, 5, 6]

$$\frac{N_0}{N} = 1 - \left(\frac{T}{T_0}\right)^3 - \frac{3\zeta(2)}{2\zeta(3)^{2/3}N^{1/3}} \left(\frac{T}{T_0}\right)^2 + \dots \quad (4)$$

The author's prediction of the quantum depletion 0.2% provides that the transition temperature should be below  $40 \sim 90\text{nK}$  for  $N = 10^5 \sim 10^6$ . Actually, the author's estimated value  $\eta_0 = 0.2\%$  is too small for the inhomogeneous system to be accepted. We obtain the fugacity  $z(T)$  first by solving Eq. (4) numerically, and then obtain the chemical potential  $\mu(T)$ . We use this  $\mu(T)$  for the estimation of  $N_0^{int}$ .

Comparing  $N_0$  and  $N_0^{int}$ , the over all behavior of the condensate fraction in Eq. (2) as a function of scattering length and temperature can be estimated by the approximation

$$N_0^{int}(a, T) \sim N_0(T)e^{-\beta U(a, N_0)}. \quad (5)$$

The ground state occupation number falls exponentially as the interaction increases.

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